

TCAD analysis of wide-spectrum waveguides in high-voltage SOI-CMOS

Satadal Dutta¹, Luis Orbe², and Jurriaan Schmitz¹.

¹ MESA+ Institute for Nanotechnology, University of Twente, ² Phoenix B.V., Enschede, The Netherlands.

Abstract—A TCAD based analysis is presented on the transmission efficiency η of silicon-on-insulator (SOI) and silicon nitride slab waveguides in a high-voltage standard SOI-CMOS technology, for the spectral range of 480 nm - 1300 nm, and isotropic optical excitation via monolithic Si-based LEDs. The effects of geometry, wavelength and galvanic isolation on η are reported.

The integration of photonic functionality in CMOS is promising for high-speed data communication, and opto-electronic system-on-chip applications. For most contemporary photonic and/or opto-electronic integrated circuits [1]–[5], CMOS technology is commonly extended with a dedicated waveguide (WG) layer having a low material absorption coefficient (α) and a high refractive index (n) for photonic functionalities [1]. Some industrial CMOS technologies, however, offer built-in thin films suitable as WGs, the most common being the silicon (Si) layer in silicon-on-insulator (SOI) technology [1], [5]–[7]. Recently, a monolithic optocoupler was realized in standard high-voltage SOI CMOS [6], which not only includes an SOI layer as a WG for infrared (IR) light, but also a thin silicon nitride (Si_3N_4) film laid atop the active Si surface, which is a potential WG for visible and IR light [7]–[12]. In addition, the optocoupler comprises an Si LED that exhibits wide-spectrum ($400 \text{ nm} < \lambda < 1300 \text{ nm}$) and isotropic electroluminescence (EL). An Si photodiode (PD) detects light laterally. Such features of the LED, combined with the inherent off-axis alignment of the WGs w.r.t. the LED and the PD, makes the optical transmission efficiency $\eta(\lambda)$ rather difficult to analyze by means other than numerical TCAD simulation.

In this work we first show, using raytracing simulations in Sentaurus, the built-in WG conditions for SOI and Si_3N_4 core layers with SiO_2 cladding, leading to anisotropic transmission. Secondly, via hybrid-mode EM wave simulations, we show how $\eta(\lambda)$ is affected by geometry; namely SOI and/or nitride thickness $t_{\text{SOI}/\text{Si}_3\text{N}_4}$, link length L , and galvanic isolation in a typical SOI-based optocoupler, with on-chip isotropic optical excitation.

Fig. 1(a) shows the schematic cross-section of the key features of a typical SOI-based optocoupler; a p-n junction Si LED and PD, the shallow trench isolation (STI) of length L , the medium trench isolation (MTI), the buried oxide (BOX), the relatively thin Si_3N_4 layer, and a back-end SiO_2 layer comprising the inter-metal dielectric (IMD). The MTI column, used for galvanic isolation, is typically composed of SiO_2 enclosing a thin Si core. A light ray originating from a point in Si along the p-n junction (EL region of the LED) makes an angle ϕ w.r.t. the positive x -axis ($-90^\circ < \phi < 90^\circ$), where ϕ determines the photon trajectory in the optocoupler. Waveguiding

via the Si_3N_4 WG requires two necessary conditions. Firstly, $n_{\text{Si}}(\lambda) > n_{\text{Si}_3\text{N}_4}(\lambda) > n_{\text{SiO}_2}(\lambda)$, and secondly, $\cos^{-1}(n_{\text{SiO}_2}(\lambda)/n_{\text{Si}}(\lambda)) > \phi > \cos^{-1}(n_{\text{Si}_3\text{N}_4}(\lambda)/n_{\text{Si}}(\lambda))$. Waveguiding via the SOI WG (without MTI) occurs only if $\phi > -\sin^{-1}(n_{\text{SiO}_2}(\lambda)/n_{\text{Si}}(\lambda))$. If an MTI column is present in the SOI layer, then necessary WG condition is additionally constrained by $\phi > -\cos^{-1}(n_{\text{SiO}_2}(\lambda)/n_{\text{Si}}(\lambda))$. The constraints on ϕ affect the optical transmission, which is captured by extracting the TCAD simulated gain G_{opt} in the PD photo-current I_{PD} , as summarized in Fig. 1(b) at $\lambda = 1100 \text{ nm}$. This choice of λ ensures negligible material absorption in both the WGs. In addition, line-of-sight (LOS) propagation (incurs mainly Fresnel reflection losses) occurs along the x -axis via a small aperture $|\phi| < \delta$ through the STI, with δ depending on L , and t_{STI} .

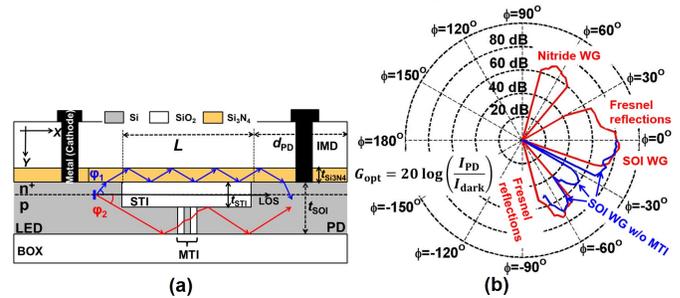


Fig. 1. (a) 2-D schematic cross-section of the SOI-based optocoupler showing relevant parameters and design features, and two example ray-traces when guided via the Si_3N_4 (blue) and the SOI WGs (red). (b) TCAD simulated G_{opt} of the PD versus ϕ at $\lambda = 1100 \text{ nm}$.

Next, the effect of geometry and λ on η for propagation via the SOI WG is studied using 2-D EM wave solver in Sentaurus. Fig. 2(a) shows the structure used as our simulation input deck. The structure is optimized to ensure a low self-absorption of light within the Si LED. A truncated plane-wave (TPW) excitation with mixed TE and TM polarization and a fixed intensity that has spatial divergence, is used to mimic our Si-embedded optical source. The TCAD simulated profiles for optical intensity and the magnetic field intensity $H(x, y)$ are shown for $\lambda = 1100 \text{ nm}$ in Figs. 2(b) and (c) respectively, showing optical confinement and guiding via the SOI layer. In Fig. 2(d), we observe that for a fixed $t_{\text{SOI}} = 1 \mu\text{m}$, and any given L , $\eta = P_{\text{out}}/P_{\text{in}}$ first increases with increasing λ due to a sharp decrease in α_{Si} , reaches a maximum at a certain λ_{peak} , and eventually decreases gradually due to a reduction in the mode-propagation efficiency for longer λ . Further, for a fixed λ , η decreases with increasing L due to absorption. For $\lambda > 1150 \text{ nm}$ (corresponds to Si band gap

of 1.12 eV), η is significantly less sensitive to L , due to negligible absorption at such long λ . In Fig. 2(e), we observe that for a fixed L , $\eta(\lambda)$ increases as t_{SOI} is increased from 1 μm to 2 μm , due to increased mode-propagation efficiency. This also explains the observation that $\eta(\lambda)$ increases till $\lambda = 1300$ nm for $t_{\text{SOI}} = 2$ μm ; λ_{peak} being shifted beyond 1300 nm. Further, the presence of the MTI, leads to a fixed additional 20 % reduction in η for any L and t_{SOI} , caused by Fresnel reflections at the Si-SiO₂ interfaces. The ripples in $\eta(\lambda)$ are likely caused by inter-modal interference.

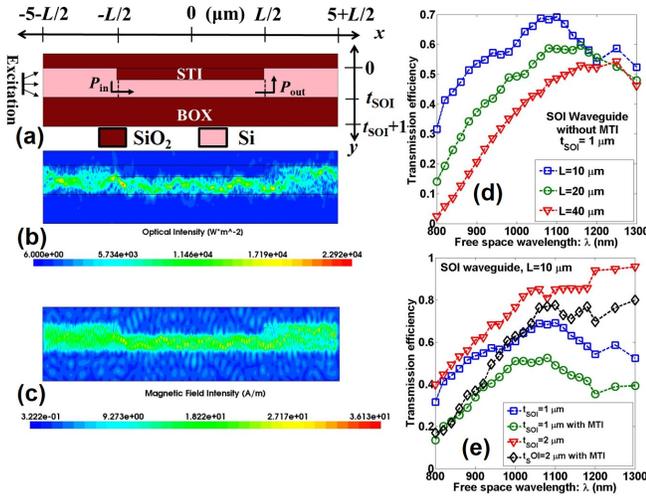


Fig. 2. (a) The SOI WG: The time-averaged and y -integrated input (P_{in}) and output (P_{out}) optical powers are evaluated at $x = -0.5L$ and $x = 0.5L$ respectively. Convolutional Perfectly Matched Layer (CPML) boundary conditions are assumed in both x and y directions, and $t_{\text{STI}} = 0.4$ μm . (b) Optical intensity profile and (c) Magnetic field intensity profile at $\lambda = 1100$ nm. (d) and (e) Simulated η for indicated values of L and t_{SOI} .

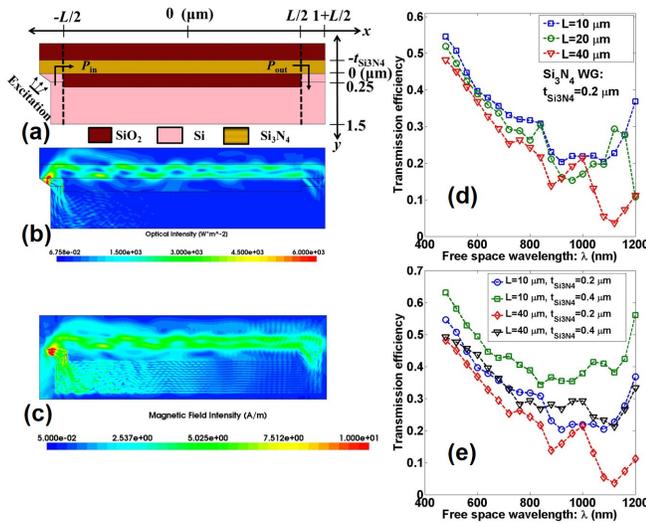


Fig. 3. (a) The Si₃N₄ WG: Time-averaged and y -integrated input (P_{in}) and output (P_{out}) optical powers are evaluated at $x = -0.5L$ and $x = 0.5L$ respectively. CPML boundary conditions are assumed for both x and y axes. (b) Optical intensity profile and (c) Magnetic field intensity profile at $\lambda = 600$ nm. (d) and (e) Simulated η for indicated values of L and $t_{\text{Si}_3\text{N}_4}$.

Fig. 3(a) shows the optimized structure used to simulate the Si₃N₄ WG. The TPW excitation (constrained by the solver to

be placed in vacuum) is given an initial orientation $\phi \approx 62^\circ$. This mimics our Si-based optical source without violating the aforesaid necessary conditions for waveguiding. Figs. 3(b) and (c) show the intensity profile and $H(x, y)$ respectively at $\lambda = 600$ nm, showing out-of-plane guiding. Light is coupled in from the Si LED into Si₃N₄, and is coupled out from Si₃N₄ to the Si PD. The Si₃N₄ core is optically much thinner than the SOI core. In Fig. 3(d), we observe that for a fixed $t_{\text{Si}_3\text{N}_4}$ and L , η falls sharply with increasing λ . A slight increase in η is observed for $\lambda \geq 1100$ nm, which can be explained by an increased edge-coupling to the underlying SOI layer, thereby augmenting P_{out} and hence η for longer λ . Further, $\eta(\lambda)$ shows a much smaller L -dependence (attenuation) than observed in the SOI case, due to the much lower $\alpha_{\text{Si}_3\text{N}_4}$. In real structures, however, a higher attenuation is expected [6] because of roughness in the core-cladding interface (ignored in our simulation). For any L , as $t_{\text{Si}_3\text{N}_4}$ is doubled from 0.2 μm to 0.4 μm , a $\sim 15\%$ increase in $\eta(\lambda)$ is observed due to increased mode-propagation efficiency; the increase being more pronounced at longer λ .

The $\eta(\lambda)$ for all the optical paths can be combined to obtain the cumulative weighted $\eta(\lambda)$ over the entire range of ϕ to mimic an isotropic optical source, as shown in Fig. 4(a). Transmission is dominated by the Si₃N₄ WG for $\lambda < 700$ nm, and by the SOI WG for $\lambda > 700$ nm. Our TCAD analysis provides some key guidelines for designers to optimize optical propagation in standard SOI CMOS technology, involving embedded wide-spectrum and isotropic light sources, as highlighted in Fig. 4(b).

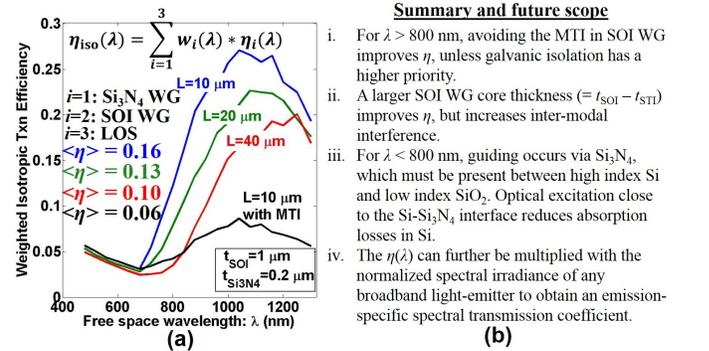


Fig. 4. (a) ϕ -weighted cumulative transmission efficiencies for indicated geometrical parameters assuming isotropic EL combining the three optical paths: Si₃N₄ WG, SOI WG, and LOS. The weights $w_i(\lambda)$ are calculated using the allowed range of ϕ for each path, such that $\sum w_i(\lambda) = 1$. The mean spectral efficiency $\langle \eta \rangle$ is indicated for each case. (b) Summary of the key take-away messages of our work and its potential scope of application.

REFERENCES

[1] C. Sun et al., Nature, vol. 528, pp. 534-538, 2015.
 [2] M. du Plessis et al., IEEE J. Sel. Topics Quantum Electron., vol. 8, no. 6, pp. 1412-1419, 2002.
 [3] B. Huang et al., Opt. Commun., vol. 284, no. 16-17, pp. 3924-3927, 2011.
 [4] A.Z. Subramanian et al., Proc. 45th ESSDERC, pp. 138-141, 2015.
 [5] N. Rouger et al., Proc. 28th ISPSD, pp. 427-430, 2016.
 [6] S. Dutta et al., Opt. Express, vol. 25, no.4, pp. 5440-5456, 2017.
 [7] R. Baets et al., in proc. OFC 2016, p. Th3J.1.
 [8] A.Z. Subramanian et al., IEEE Photonics J., vol. 5, no. 6, 2013.
 [9] M. Melchiorri et al., Appl. Phys. Lett., vol. 86, 121111, 2005.
 [10] H. Yamada, Piers Online, vol. 6, no. 2, 2010.
 [11] G. Li et al, Opt. Express, vol. 20, no. 11, pp. 12035-12039, 2012.
 [12] F. Grillot et al., J. Lightwave Tech., vol. 24, no. 2, 2006.